ERRATUM: EULERIAN DYNAMICS WITH A COMMUTATOR FORCING

ROMAN SHVYDKOY AND EITAN TADMOR

The publication [1] has a minor gap in the argument presented in Section 6.2 where the authors establish control over the first derivatives of density and momentum. Specifically, the bound on $\Lambda\rho$ used in the momentum equation involves term $\sqrt{D\rho'(x)}$, which propagates into formula (6.21). At that point the authors combined (6.21) with (6.19) to get rid of the *D*-term. The mistakes presents in the fact that the point *x* at which the *D*-term is evaluated in 6.19 is different from the point *x* at which it is evaluated in 6.21. Hence the values may be different.

To avoid using combination of 6.19 and 6.21 we argue as follows. We produce a uniform bound on $|\rho''|_2$ on the time interval in question. This uniform bound, by Sobolev embedding, implies that $\rho' \in C^{\frac{1}{2}}$ uniformly. Then the trivial bound

$$|\Lambda\rho|_{\infty} \leqslant |\rho'|_{C^{1/2}},$$

implies uniform control over $\Lambda \rho$. Hence it is not necessary to resort to 6.19 to contain $\Lambda \rho$, and the rest of the estimates on m' follow as documented in [1].

To achieve uniform bound on $|\rho''|_2$ we differentiate the density equation twice:

$$\partial_t \rho'' + u\rho''' + u'\rho'' + e''\rho + 3e'\rho' + 2e\rho'' = -2\rho''\Lambda\rho - 3\rho'\Lambda\rho' - \rho\Lambda\rho''.$$

Using that $u' = e + \Lambda \rho$, we obtain

$$\partial_t \rho'' + u \rho''' + e'' \rho + 3e' \rho' + 3e \rho'' = -3\rho'' \Lambda \rho - 3\rho' \Lambda \rho' - \rho \Lambda \rho''.$$

At this point we know that $|e^{(k)}| \leq \rho^{(k)}$, and we have uniform bounds on ρ, ρ' . So, testing with ρ'' , integrating by parts in $u\rho'''\rho''$ term, and using the *e* quantity again, we obtain

$$\partial_t |\rho''|_2^2 \lesssim |\rho''|_2 + |\rho''|_2^2 + |\Lambda\rho|_{\infty} |\rho''|_2^2 + |\rho''|_2 |\Lambda\rho'|_2 - \int_{\mathbb{T}} \rho \rho'' \Lambda \rho'' dx.$$

Using that $|\Lambda \rho'|_2 \lesssim |\rho''|_2$, and log-Sobolev inequality

$$|\Lambda\rho|_{\infty} \leqslant |\rho'|_{\infty} (1 + \log_+ |\rho''|_2) \lesssim 1 + \log_+ |\rho''|_2,$$

we further obtain

$$\partial_t |\rho''|_2^2 \lesssim C + |\rho''|_2^2 (1 + \log_+ |\rho''|_2) - \int_{\mathbb{T}} \rho \rho'' \Lambda \rho'' dx.$$

Using symmetrization in the remaining dissipation term we have

(0.1)
$$-\int_{\mathbb{T}}\rho\rho''\Lambda\rho''dx = -\int_{\mathbb{T}}\rho D\rho''dx + R,$$

where

$$R = \int_{\mathbb{T}} \rho''(x) \int_{\mathbb{T}} \frac{(\rho(x) - \rho(y))(\rho''(x) - \rho''(y))}{|x - y|^2} dy dx$$

Date: October 21, 2018.

Using the bound $|\rho'| < C$ we further conclude

$$|R| \lesssim \int_{\mathbb{T}} |\rho''(x)| \int_{\mathbb{T}} \frac{|\rho''(x) - \rho''(y)|}{|x - y|} dy dx \leqslant \int_{\mathbb{T}} |\rho''(x)| \sqrt{D\rho''} dx \leqslant |\rho''|_2 \sqrt{\int_{\mathbb{T}} D\rho'' dx}.$$

By Young, the latter is bounded by

$$|R| \leqslant \varepsilon \int_{\mathbb{T}} D\rho'' dx + C_{\varepsilon} |\rho''|_2^2$$

where ε is smaller than the lower bound on the density on the given time interval. This gets the *D*-term absorbed into dissipation term in (0.1). We thus arrive at

$$\partial_t |\rho''|_2^2 \lesssim C + |\rho''|_2^2 (1 + \log_+ |\rho''|_2).$$

The result follows by integration.

Since in the estimates above we relied on second order a priori bound $|e''| \leq |\rho''|$ it is necessary to raise the regularity class from H^3 as in [1] to H^4 so that the local transport equation for e'' can be solved classically. The idea to avoid using higher order a priori bounds $|e^{(k)}| \leq |\rho^{(k)}|$ is to abandon the use of momentum equation for m, where e quantity is explicitly present, and instead come back to the *u*-equation. This was performed in [1] up to the order 3 space H^4 , and the argument is entirely similar going one more derivative up to H^4 . We therefore state our final result as follows.

Theorem 0.1. Consider the system of equations (1.1), [1], with $1 \leq \alpha < 2$ subject to initial data $(u_0, \rho_0) \in H^4(\mathbb{T}^1) \times H^{3+\alpha}(\mathbb{T}^1)$. Then the system admits a global solution in the same class.

References

 R. Shvydkoy and E. Tadmor, Eulerian dynamics with a commutator forcing, Trans. Math. and Appl. 1(1) (2017) 1-26.

Department of Mathematics, Statistics, and Computer Science, M/C 249,, University of Illinois, Chicago, IL 60607, USA

Email address: shvydkoy@uic.edu

DEPARTMENT OF MATHEMATICS, CENTER FOR SCIENTIFIC COMPUTATION AND MATHEMATICAL MOD-ELING (CSCAMM), AND INSTITUTE FOR PHYSICAL SCIENCES & TECHNOLOGY (IPST), UNIVERSITY OF MARYLAND, COLLEGE PARK

CURRENT ADDRESS: INSTITUTE FOR THEORETICAL STUDIES (ITS), ETH, CLAUSIUSSTRASSE 47, CH-8092 ZURICH, SWITZERLAND

Email address: tadmor@cscamm.umd.edu